# Structural biomathematics: an overview of molecular simulations and protein structure prediction 

## Bernat Anton




Figure: Parc de Recerca Biomèdica de Barcelona (PRBB).

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(1) A Glance at Structural Biology
(2) Molecular Simulations

3 Direct-Coupling Analysis for Prediction of Protein Folding

## (9) A Glance at Structural Biology

## (2) Molecular Simulations

3 Direct-Coupling Analysis for Prediction of Protein Folding

All the biological information of the human body is encoded in our DNA. Human Genome Project: Sequentiation of the whole human genome completed on 2001, by Francis Collins (Public Project) \& Craig Venter (Celera Genomics).

- About 3 billion base pairs (A, C, T and G).
- Estimation of 30000 genes (around 3000bp per gene).
- Less than $2 \%$ of the genome codes for proteins.
- Unknown function for over the half of the discovered genes!

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ATGAAAAAGACAGCTATCGCGATTGCAGTGGCACTGGCTGGTTTCGCTACCGTGGCC CAGGCGGCCTCTGAGGGAAACAGTGACTGCTACTTTGGGAATGGGTCAGCCTACCG TGGCACGCACAGCCTCACCGAGTCGGGTGCCTCCTGCCTCCCGTGGAATTCCATGAT CCTGATAGGCAAGGTTTACACAGCACAGAACCCCAGTGCCCAGGCACTGGGCCTGG GCAAACATAATTACTGCCGGAATCCTGATGGGGATGCCAAGCCCTGGTGCCACGTG CTGAAGAACCGCAGGCTGACGTGGGAGTACTGTGATGTGCCCTCCTGCTCCACCTGC GGCCTGAGACAGTACAGCCAGCCTCAGTTTCGCATCAAAGGAGGGCTCTTCGCCGA CATCGCCTCCCACCCCTGGCAGGCTGCCATCTTTGCCAAGCACAGGAGGTCGCCCGG AGAGCGGTTCCTGTGCGGGGGCATACTCATCAGCTCCTGCTGGATTCTCTCTGCCGC CCACTGCTTCCAGGAGAGGTTTCCGCCCCACCACCTGACGGTGATCTTGGGCAGAAC ATACCGGGTGGTCCCTGGCGAGGAGGAGCAGAAATTTGAAGTCGAAAAATACATTG TCCATAAGGAATTCGATGATGACACTTACGACAATGACATTGCGCTGCTGCAGCTGA AATCGGATTCGTCCCGCTGTGCCCAGGAGAGCAGCGTGGTCCGCACTGTGTGCCTTC CCCCGGCGGACCTGCAGCTGCCGGACTGGACGGAGTGTGAGCTCTCCGGCTACGGC AAGCATGAGGCCTTGTCTCCTTTCTATTCGGAGCGGCTGAAGGAGGCTCATGTCAGA CTGTACCCATCCAGCCGCTGCACATCACAACATTTACTTAACAGAACAGTCACCGAC AACATGCTGTGTGCTGGAGACACTCGGAGCGGCGGGCCCCAGGCAAACTTGCACGA CGCCTGCCAGGGCGATTCGGGAGGCCCCCTGGTGTGTCTGAACGATGGCCGCATGA CTTTGGTGGGCATCATCAGCTGGGGCCTGGGCTGTGGACAGAAGGATGTCCCGGGT GTGTACACAAAGGTTACCAACTACCTAGACTGGATTCGTGACAACATGCGACCG (SEO ID NO:2)


## DNA $\xrightarrow{\text { Transcription }}$ RNA $\xrightarrow{\text { Translation }}$ Protein

| Standard genetic code |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1st |  |  |  |  | ase |  |  |  | 3rd |
| base |  | U |  | C |  | A |  | G | base |
|  | UUU | nine | UCU |  | UAU |  | UGU |  | U |
| U | UUC |  | UCC |  | UAC |  | UGC |  | C |
| U | UUA | (Leu/L) Leucine | UCA |  | UAA | Stop (Ochre) | UGA | Stop (Opal) | A |
|  | UUG |  | UCG |  | UAG | Stop (Amber) | UGG | (Trp/W) Tryptophan | G |
| C | CUU |  | CCU | (Pro/P) Proline | CAU | (His/H) Histidine | CGU | (Arg/R) Arginine | U |
|  | CUC |  | CCC |  | CAC |  | CGC |  | C |
|  | CUA |  | CCA |  | CAA | (Gln/Q) Glutamine | CGA |  | A |
|  | CUG |  | CCG |  | CAG |  | CGG |  | G |
| A | AUU | (lle/l) Isoleucine | ACU | (Thr/T) Threonine | AAU | (Asn/N) Asparagine | AGU | (Ser/S) Serine | U |
|  | AUC |  | ACC |  | AAC |  | AGC |  | C |
|  | AUA |  | ACA |  | AAA | (Lys/K) Lysine | AGA | (Arg/R) Arginine | A |
|  | $A \cup G{ }^{[A]}$ | (Met/M) Methionine | ACG |  | AAG |  | AGG |  | G |
| G | GUU | $(\mathrm{Val} / \mathrm{V})$ Valine | GCU | (Ala/A) Alanine | GAU | (Asp/D) Aspartic acid | GGU | (Gly/G) Glycine | U |
|  | GUC |  | GCC |  | GAC |  | GGC |  | C |
|  | GUA |  | GCA |  | GAA | (Glu/E) Glutamic acid | GGA |  | A |
|  | GUG |  | GCG |  | GAG |  | GGG |  | G |

[^0]

Protein structure $\xrightarrow{? ? ?}$ Protein function $\longrightarrow$ Gene function


- Primary: Amino acid linear sequence.
- Secondary: $\alpha$-helices and $\beta$-strands.
- Tertiary / Domains: Functionally independent part of the sequence.
- Quaternary: Multi-subunit complex of domains or proteins.

[^1]
## Main question:

How can we find the structure of a given protein?

- Crystallography.
- Nuclear magnetic resonance spectroscopy.
- Molecular simulation.
- Prediction of structure (structural biology).

NOT AN EASY TASK!

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## (1) A Glance at Structural Biology

(2) Molecular Simulations

## 3 Direct-Coupling Analysis for Prediction of Protein Folding



3
${ }^{3}$ Both images were obtained using VMD software


Lysine


Internal (mechanical) energy of the system

And these are not the only forces and energies implied in a molecular simulation!

PLC- $\beta 2$ simulation
This simulation lasts around $20 n s$, with timesteps of $4 f s^{4}$, using the ACEMD software with the AMBER forcefield. The simulation has been visualized using VMD software.
The protein has 708 amino acids, for a total of around 150000 atoms in the simulation (counting water and lipid molecules).

In the simulation can be observed the folding of the $X / Y$ linker in order to cover the hydrophobic active site of the protein.

$$
{ }^{4} 1 \mathrm{~ns}=10^{-9} \text { seconds, } 1 \mathrm{fs}=10^{-15} \text { seconds }
$$

## Afinsen's Dogma

The native structure of a protein is unique and is determined only by it's amino acid sequence. The folding to its native state is almost spontaneous.

## Levinthal's Paradox

Due to the huge number of degrees of freedom in an unfolded protein, the number of possible conformations is astronomically large.

Then... how can proteins fold?

- Partially folded transition states.
- Funnel-like energy landscapes.


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- ...?


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## (2) <br> Molecular Simulations

(3) Direct-Coupling Analysis for Prediction of Protein Folding

Let $X, Y$ be two (discrete) random variables.

- The (self-)information of $X$ is $I(X)=-\log (P(X))$.
- The entropy of $X$ is the measure of uncertainty associated with $X: S(X)=E(I(X))$.
- The mutual information of $X$ and $Y$ (also called Kullback-Leibler divergence) is

$$
M I(X ; Y)=\sum_{x \in X} \sum_{y \in Y} P(x, y) \log \left(\frac{P(x, y)}{P(x) P(y)}\right)
$$

## Maximum Entropy Principle

Given a proposition that expresses testable information, the probability distribution that best represents the current state of knowledge is the one with largest entropy.


Figure: Multiple Sequence Alignment (MSA) for aaTHEP1.

From the previous MSA let's define:
$f_{i}(A)=$ 'frequency of apparitions of aa $A$ in the position $i$ of the MSA'
$f_{i, j}(A, B)=$ 'frequency of simultaneous apparitions of aa $A$ and $B$ in respective positions $i$ and $j$ of the MSA'

$$
M I_{i, j}=\sum_{A, B} f_{i, j}(A, B) \ln \left(\frac{f_{i, j}(A, B)}{f_{i}(A) f_{j}(B)}\right)
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> Be careful!
> By definition, this mutual information of these frequencies is local in the amino acid chain, thus is noised by transitivity of correlations.

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## We want

$$
P\left(A_{1}, \ldots, A_{L}\right)
$$

a general model for the probability of a particular amino acid sequence $A_{1} \ldots A_{L}$ to be member of the iso-structural family under consideration, and such that

$$
\begin{aligned}
P_{i}(A) & \approx f_{i}(A), \\
P_{i, j}(A, B) & \approx f_{i, j}(A, B),
\end{aligned}
$$

where

$$
\begin{aligned}
P_{i}(A) & =\sum_{A_{k} \neq A} P\left(A_{1}, \ldots, A_{L}\right), \\
P_{i, j}(A, B) & :=\sum_{A_{k} \neq A, B} P\left(A_{1}, \ldots, A_{L}\right) .
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Many distributions satisfying this: Maximum Entropy Principle!!!!

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## Optimization problem:

maximize $S=-\sum_{A_{i} \mid i=1, \ldots, L} P\left(A_{1}, \ldots, A_{L}\right) \ln P\left(A_{1}, \ldots, P_{L}\right)$
subject to

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\begin{aligned}
P_{i, j}(A, B) & =f_{i, j}(A, B) \\
P_{i}(A) & =f_{i}(A)
\end{aligned}
$$

## Solution: disordered Q-state Potts model


where:

- the parameters $e_{i, j}\left(A_{i}, A_{j}\right), h_{i}\left(A_{i}\right)$ are the Lagrange multipliers of the system,
- $\mathcal{Z}$ is the normalization constant (partition function).

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Solution: disordered Q-state Potts model

$$
P\left(A_{1}, \ldots, A_{L}\right)=\frac{1}{\mathcal{Z}} \exp \left\{\sum_{1 \leq i<j \leq L} e_{i, j}\left(A_{i}, A_{j}\right)+\sum_{1 \leq i \leq L} h_{i}\left(A_{i}\right)\right\}
$$

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Geometrically, this probability distribution is given by the Boltzmann-Gibbs distribution:

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P\left(A_{1}, \ldots, A_{L}\right)=\frac{1}{\mathcal{Z}} e^{-\mathcal{H}\left(A_{1}, \ldots, A_{L}\right)}
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but the direct computation is computationally prohibitive.

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Formally, the marginals of this distribution are obtained from

$$
\begin{aligned}
\frac{\partial \ln \mathcal{Z}}{\partial h_{i}(A)} & =-P_{i}(A) \\
\frac{\partial^{2} \ln \mathcal{Z}}{\partial h_{i}(A) \partial h_{j}(B)} & =-P_{i, j}(A, B)+P_{i}(A) P_{j}(B)
\end{aligned}
$$

but the direct computation is computationally prohibitive.

## The Lagrange multipliers can be obtained using Mean Field Aproximation technique ${ }^{5}$ :

- Introduce a new parameter $\alpha$ in the partition function (via the disturbed Hamiltonian):

$$
\mathcal{H}(\alpha)=\sum_{i=1, \ldots, L} \exp \left\{\alpha \sum_{1 \leq i<j \leq L} e_{i, j}\left(A_{i}, A_{j}\right)+\sum_{1 \leq i \leq L} h_{i}\left(A_{i}\right)\right\}
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- Consider the Legendre transform of the Gibbs free energy $\mathcal{F}=-\ln \mathcal{Z}(\alpha)$ (Gibbs potential):


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$$

- Consider the Legendre transform of the Gibbs free energy $\mathcal{F}=-\ln \mathcal{Z}(\alpha)$ (Gibbs potential):

$$
\mathcal{G}(\alpha)=\ln \mathcal{Z}(\alpha)-\sum_{i=1, \ldots, L} \sum_{A} h_{i}(A) P_{i}(A) .
$$

[^3]- Considering the empirical connected correlation matrix:

$$
C_{i, j}(A, B)=f_{i, j}(A, B)-f_{i}(A) f_{j}(B) .
$$

As a consecuence of the functional form of the Legendre transform

$$
\begin{aligned}
h_{i}(A) & =\frac{\partial \mathcal{G}(\alpha)}{\partial P_{i}(A)} \\
\left(C^{-1}\right)_{i,}(A, B) & =\frac{\partial h_{i}(A)}{\partial P_{j}(B)}=\frac{\partial^{2} \mathcal{G}(\alpha)}{\partial P_{i}(A) \partial P_{j}(B)}
\end{aligned}
$$

- Expand the Gibbs potential up to first order Taylor expansion around $\alpha=0$ :

$$
\mathcal{G}(\alpha) \approx \mathcal{G}(0)+\alpha{\frac{\partial \mathcal{G}(\alpha)}{\partial \alpha}{ }_{\mid \alpha=0} .}
$$

A computation over the two terms of the Taylor expansion of $\mathcal{G}$ leads us to an expression which is easily derivable. First and second derivatives with respect the marginal distributions $P_{i}(A)$ provide self-consistent equations for the local fields, from which we obtain

$$
\left(C^{-1}\right)_{i, j}(A, B)_{\mid \alpha=0}=-e_{i, j}(A, B), \text { for } i \neq j
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Finally, the parameters $h_{i}$ can be estimated imposing empirical single-site frequency counts as marginal distributions and considering gauge conditions:


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$$
f_{i}(A)=\sum_{B} P_{i, j}(A, B) .
$$

This leads us to the effective pair probabilities

$$
P_{i, j}^{\text {Dir }}(A, B)=\frac{1}{\mathcal{Z}_{i, j}} \exp \left\{e_{i, j}(A, B)+\tilde{h}_{i}(A)+\tilde{h}_{j}(B)\right\} .
$$

From which we can define its Kullback-Leibler divergence, that will be called Direct Information:

$$
D I_{i, j}:=\sum_{A, B} P_{i, j}^{D i r}(A, B) \ln \left(\frac{P_{i, j}^{D i r}(A, B)}{f_{i}(A) f_{j}(B)}\right) .
$$

## For each pair of positions in the sequence, we "know" if they are spatially close.

## And now... what?

- Depending on the previous, not an unique folding for the protein is possible. We must remove knotted structures (Alexander polynomial or Heegaard Floer homology).
- A scoring method over the resulting foldings must be defined, in order to decide which one of the structures is better.
- A short simulation of the system may be run in order to optimize the energies of the folding.

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Figure: Chewbacca mounted on a squirrel wants to thank you for your assistance!


[^0]:    ${ }^{1}$ Table taken from Wikipedia.

[^1]:    ${ }^{2}$ Figure taken from C.Branden \& J.Tooze, Introduction to Protein Structure.

[^2]:    ${ }^{5}$ Plefka, T., Convergence condition of the TAP equation for the infinite-ranged/Ising spin glass model 1982)

[^3]:    ${ }^{5}$ Plefka, T., Convergence condition of the TAP equation for the infinite-ranged/lsing spin glass mode/ 1982 )

